

Free-boundary problems — Exercises 1

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1. (Work through the nondimensionalisation from the lecture)

The two-phase Stefan problem on the rectangular domain $[0, d_1] \times [0, d_2]$ is

$$\begin{aligned} \rho c_p T_t^l &= k \nabla^2 T^l && \text{in } \Omega_l \\ \rho c_p T_t^s &= k \nabla^2 T^s && \text{in } \Omega_s \\ T^l &= T^{\text{hot}} && \text{on } x = 0, \\ \nabla T^{l \text{ or } s} \cdot \mathbf{n} &= 0 && \text{on all other fixed boundaries,} \\ T^l &= T^s = T_m && \text{on } \Gamma, \\ \rho L V &= k \nabla T^s \cdot \mathbf{n} - k \nabla T^l \cdot \mathbf{n} && \text{on } \Gamma, \end{aligned}$$

where V is the normal velocity to the free boundary Γ .

Nondimensionalise this problem using the scalings (change of variables)

$$\mathbf{x} = d_1 \hat{\mathbf{x}}, \quad t = \frac{\rho L d_1^2}{k(T^{\text{hot}} - T_m)} \hat{t}, \quad T^{l \text{ or } s} = T_m + (T^{\text{hot}} - T_m) \hat{T}^{l \text{ or } s},$$

and with $V = (k(T^{\text{hot}} - T_m)/(\rho L d_1)) \hat{V}$, and show that the dimensionless problem depends on just two dimensionless groups:

$$\text{St} = \frac{c_p(T^{\text{hot}} - T_m)}{L} \quad \text{and} \quad a = \frac{d_2}{d_1}.$$

2. (A different freezing problem)

Consider the 1D Stefan problem in the limit of $\text{St} \ll 1$ (like the lecture) but where we cool the solid rather than heating the liquid, so that we have freezing:

$$\begin{aligned} 0 &= T_{xx}^l && \text{for } x \in (0, d(t)), \\ 0 &= T_{xx}^s && \text{for } x \in (d(t), 1), \\ T^s &= -1 && \text{at } x = 1, \\ T_x^l &= 0 && \text{at } x = 0, \\ T^l = T^s = 0 &\quad \text{and} \quad d_t = T_x^s - T_x^l && \text{at } x = d(t). \end{aligned}$$

Solve this free-boundary problem to find the temperatures $T^s(x, t)$ and $T^l(x, t)$, and the position of the freezing front $x = d(t)$. Sketch d against t and the temperature profile T against x at 3 different times.

3. (Is freezing always unstable?)

Consider a similar freezing problem to q2, but in 2D and on an infinite domain:

$$\begin{aligned} 0 &= T_{xx}^l + T_{yy}^l && \text{for } x < \gamma(y, t), \\ 0 &= T_{xx}^s + T_{yy}^s && \text{for } x > \gamma(y, t), \\ T_x^s &\rightarrow -v && \text{as } x \rightarrow \infty, \\ T_x^l &\rightarrow 0 && \text{as } x \rightarrow -\infty, \\ T^l = T^s = 0 &\quad \text{and} \quad V = \nabla T^s \cdot \mathbf{n} - \nabla T^l \cdot \mathbf{n} && \text{at } x = \gamma(y, t). \end{aligned}$$

Show that a base-state solution is

$$T_B^l = 0, \quad T_B^s = -v(x + vt), \quad \gamma_B = -vt.$$

(if we assume $\gamma = 0$ at $t = 0$).

Perform a linear stability analysis to determine the stability of this base-state solution:

- (a) Make a small perturbation about the base state:

$$T^l = T_B^l + \theta^l, \quad T^s = T_B^s + \theta^s, \quad \gamma(y, t) = -vt + \xi(y, t),$$

where all of θ^l, θ^s, ξ are small. Substitute in and linearise to find a linear system (PDEs and BCs) for θ^l, θ^s and ξ .

- (b) Look for Fourier mode solutions of the form

$$\theta^l = A(x)e^{iky+\sigma t}, \quad \theta^s = B(x)e^{iky+\sigma t}, \quad \xi = e^{iky+\sigma t},$$

and solve for $A(x)$ and $B(x)$. Hence show that the growth rate σ of the perturbation is

$$\sigma = -vk.$$

- (c) Is this freezing problem stable or unstable?

4. (What if St is not small?)

So far we've used $St \ll 1$ to simplify the PDEs, but it's possible to solve the full 1D problem when St is not negligible, using a *similarity solution*. (You may have seen these in your degree. If not, here's a first example, and we'll look at another in the lecture later...)

Consider the 1D Stefan problem

$$\begin{aligned} St T_t^l &= T_{xx}^l && \text{in } x \in (0, d(t)) \\ St T_t^s &= T_{xx}^s && \text{in } x \in (d(t), 1) \\ T^l &= 1 && \text{on } x = 0, \\ T_x^s &= 0 && \text{on } x = 1, \\ T^l &= T^s = 0 && \text{on } x = d(t), \\ d_t &= T_x^s - T_x^l && \text{on } x = d(t), \\ d &= 0 \quad \text{and} \quad T^s = T^l = 0 && \text{at } t = 0. \end{aligned}$$

- (a) Show that $T^s = 0$ everywhere satisfies the PDE for T^s and the Dirichlet boundary conditions. Hence reduce the problem to just a system for T^l and d .
- (b) We'll now take advantage of the inherent symmetry of the problem to reduce the PDE system to an ODE system. Look for a similarity solution of the form:

$$T^l = f(\eta), \quad d(t) = \zeta\sqrt{t} \quad \text{where} \quad \eta = \frac{x}{\sqrt{t}} \quad \text{and } \zeta \text{ is a constant,} \quad (1)$$

ie: using the chain rule, show that the change of variables means derivatives are replaced by

$$\frac{\partial}{\partial x} = \frac{1}{\sqrt{t}} \frac{d}{d\eta}, \quad \frac{\partial}{\partial t} = -\frac{\eta}{2t} \frac{d}{d\eta}.$$

Thus show that the PDE system reduces to the ODE system

$$\begin{aligned} \frac{d^2 f}{d\eta^2} + St \frac{\eta}{2} \frac{df}{d\eta} &= 0, && \text{for } \eta \in (0, \zeta), \\ f(0) &= 1, \\ f(\zeta) &= 0, \\ \frac{df}{d\eta}(\zeta) &= -\frac{\zeta}{2}. \end{aligned}$$

- (c) Use an integrating factor to take a first integral of the ODE for f . Hence find an integral expression for $f(\eta)$. Apply the boundary conditions to fix the integration constants and to determine an integral equation for ζ (this must be solved numerically).
- (d) Look up “error function, erf” online, and use this to simplify your expression for $f(\eta)$.
- (e) Convert back to the original variables, to find $T^l(x, t)$ and $d(t)$, in terms of ζ .
- (f) Show that in the limit $St \rightarrow 0$ we regain the linear-in- x solution from lectures.
- (g) Use an online solver, or a programming language like python or matlab to solve the equation for ζ numerically.

- (h) Sketch $f(\eta)$, perhaps using an online plotting tool or your favourite programming language to help.
- (i) Use your sketch of f to sketch $T^l(x, t)$.
- (j) How could you have worked out the form (1) for the similarity solution if I hadn't told you? Why is there a similarity solution for this problem? Can you come up with other boundary conditions for a freezing or melting Stefan problem for which a similarity solution also works? Can you find choices of boundary conditions for which it doesn't work?